5.2 Caloric properties of ideal-gas mixtures

To calculate the chemical potential μ_i of one component *i* we combine Eq. (3.78) and $\mu = \Delta G_m$

$$\mu_i = \mu_i^0 + RT \ln \frac{p_i}{p^0}$$
 with $p_i = x_i p_{tot} \Rightarrow \mu_i = \mu_i^0 + RT \ln x_i + RT \ln \frac{p_{tot}}{p^0}$ (5.1)

This is Dalton's law. Normally p_{tot} and p^0 (the standard pressure) are selected 1 bar, thus, the second ln-term is vanishing.

For an extensive property Φ the difference by mixing compared to the pure components is

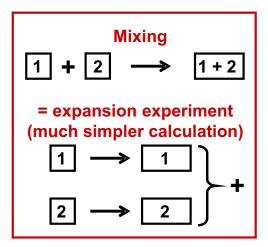
$$\Delta_{mix}\Phi^{id} = \Phi^{id}(mixture) - \sum_{i} n_i \Phi_m^{id}(pure)$$
(5.2)

Since no interaction between the ideal gas molecules exist we find

$$\Delta_{mix}H^{id} = \Delta_{mix}V^{id} = 0 \tag{5.3}$$

To calculate the effect of mixing on the entropy we apply the scheme of Fig. 5.1. Since S is a state function we can replace the calculation of mixing by expansion performed separately for each component. Since this ideal mixing is isothermal (0 = dU = T dS - p dV)

$$dS = \frac{p}{T}dV = \frac{n R T}{VT}dV = n R d \ln V$$



(5.4) Figure 5.1: Scheme for mixing of two ideal gases.

$$\Delta_{mix} S_1^{id} = n_1 R \ln \frac{V_1 + V_2}{V_1}, \qquad \Delta_{mix} S_2^{id} = n_2 R \ln \frac{V_1 + V_2}{V_2}$$

$$\Rightarrow \Delta_{mix} S_1^{id} = -n_1 R \ln x_1, \qquad \Delta_{mix} S_2^{id} = -n_2 R \ln x_2$$

$$\Rightarrow \Delta_{mix} S^{id} = \Delta_{mix} S_1^{id} + \Delta_{mix} S_2^{id} = -\sum_i n_i R \ln x_i$$
(5.5)

Thus

so

$$\Delta_{mix}G^{id} = \sum_{i} n_i R T \ln x_i \quad \text{i.e.} \quad \Delta_{mix}G^{id}_m = -T \ \Delta_{mix}S^{id}_m = \sum_{i} x_i R T \ln x_i \tag{5.6}$$

Since $\Delta_{mix}S > 0$ and $\Delta_{mix}G < 0$ mixing of gases is SPONTANEOUS, i.e. an irreversible process, even for ideal gases for which no energy contributions exist.