5.4 Current flow through a non degenerated semiconductor

We investigate the charge flow in an electrical field for a non degenerated semiconductor. Combining the equations (5.18) and (5.21) we get

$$\vec{j} = -\frac{q}{4\pi^3} \int_{V_b} \tau(\vec{k}) \vec{v}(\vec{k}) \frac{\partial f_0}{\partial E} \left\langle \left(q\vec{E} - \vec{\nabla}_r \mu - (E - \mu) \vec{\nabla}_r \ln(T) \right), \vec{v} \right\rangle d^3k \qquad . \tag{5.24}$$

We will neglect gradients in the temperature, so $(E - \mu)\vec{\nabla}_r \ln(T)$ vanishes. Since the semiconductor is not degenerated we can simplify the Fermi statistics by the Boltzmann statistics, i.e.

$$\frac{\partial f_0}{\partial E} = -\frac{f_0}{kT} \qquad . ag{5.25}$$

We take the band energies of the free electron gas

$$E = E_0 + \frac{m^*}{2} \left| \vec{v}(\vec{k}) \right|^2$$
, so $v_i = \frac{\hbar}{m^*} k_i$. (5.26)

In addition we assume τ to be independent of \vec{k} . Summing up all approximations we find for the particle current

$$\vec{j} = -\frac{q\tau\hbar^2}{4\pi^3 k T(m^*)^2} \left[e\vec{E} - \vec{\nabla}\mu \right] \tilde{M} \qquad , \tag{5.27}$$

and \tilde{M} is a matrix with the components

$$\tilde{M}_{ij} = \int_{V_k} k_i k_j f_0(E(k)) d^3k = \int_k dk k^4 f_0(E(k)) \int_{S_k} d\Omega \frac{k_i k_j}{k^2} \qquad .$$
 (5.28)

Using $k^2 = k_x^2 + k_y^2 + k_z^2$ and integrating over the surface of a sphere we get

$$\int_{S_k} d\Omega \frac{k_i k_j}{k^2} = \frac{4\pi}{3} \delta_{ij} \qquad . \tag{5.29}$$

Using

$$dE = \frac{\hbar^2}{m^*} k dk \qquad , \tag{5.30}$$

the current density is written as

$$\vec{j} = -\frac{q\tau}{3\pi^2 k T m^*} \left[q\vec{E} - \vec{\nabla}\mu \right] \left(\frac{2m^*}{\hbar^2} \right) \int_{E_0}^{\infty} dE (E - E_0)^{\frac{3}{2}} \exp\left(-\frac{E - \mu}{kT} \right)$$
 (5.31)

After partial integration we get

$$\vec{j} = -\frac{q\tau}{3\pi^2 k T m^*} \left[q\vec{E} - \vec{\nabla}\mu \right] \left(\frac{2m^*}{\hbar^2} \right) \frac{-3kT}{2} \int_{E_0}^{\infty} dE (E - E_0)^{\frac{1}{2}} \exp\left(-\frac{E - \mu}{kT} \right)$$
 (5.32)

Taking into account the density of state of free electrons

$$D(E) = \frac{1}{2\pi^2} \left(\frac{2m^*}{\hbar^2}\right)^{\frac{3}{2}} (E - E_0)^{\frac{1}{2}} , \qquad (5.33)$$

we finally get

$$\vec{j} = \frac{q\tau}{m^*} \left[q\vec{E} - \vec{\nabla}\mu \right] \int_{E_0}^{\infty} dE D(E) f_0(E) = \frac{q\tau n_e}{m^*} \left[q\vec{E} - \vec{\nabla}\mu \right]$$
 (5.34)