Reciprocal Lattice

Basics

Geometric Definition

The reciprocal lattice is fundamental for all diffraction effects and other processes in a crystal lattice where momentum is transferred.

The *reciprocal lattice* of any *geometrical point lattice* has a simple geometric definition:

- It can be constructed by drawing the direct lattice, picking three sets of lattice planes **(hi , ki , li)** (**i=1,2,3**) that are not coplanar, and by constructing three vectors *gh,k,l* which are perpendicular to the respective lattice planes and with a length (measured in cm–1) that is given by **|***g***|=2π/***d***h,k,l**, with *d***h,k,l**=distance between the lattice planes **(h,k,l)**.
- The three vectors thus obtained, if reduced to the three shortest ones possible (take three lattice planes with largest distance, i.e. lowest values of **(h,k,l)**) define the reciprocal lattice.

This is, of course, just a complicated way of saying:

Take the **(100)**, **(010)**, and the **(001)** planes, and use the vectors perpendicular to those planes with a length given by **2π/***d* for these **{100}** type planes as the base vectors of the reciprocal lattice.

The Reciprocal Lattice as Fourier Transform of the Regular Lattice

The reciprocal lattice, however, is best looked at as the **Fourier [transform](http://www.tf.uni-kiel.de/matwis/amat/elmat_en/kap_3/basics/b3_3_2.html)** of the regular lattice. We are showing this by constructing the Fourier transform of a real *crystal*.

- It is easier to look at a real crystal (not just a lattice) because otherwise you have to work with **δ**-functions.
- A real crystal has atoms. And atoms contain charge densities *ρ* **(***r*), or, if we start simple and one-dimensional, **ρ (***x***)**.
- Now, **ρ(** *x***)** must be periodic in *x*-direction with the lattice constant *a*:

$$
\rho(x + na) = \rho(x),
$$
 $n=0 \pm 1, \pm 2, ...$

We thus can expand **ρ (***x***)** into a [Fourier series,](http://www.tf.uni-kiel.de/matwis/amat/elmat_en/kap_3/basics/b3_3_2.html) i.e.

$$
\rho(x) = \sum_{n=0}^{\infty} \rho_n \cdot \exp \frac{i \cdot x \cdot n \cdot 2\pi}{a}
$$

The three-dimensional case, in analogy, can be written as

$$
\rho(\underline{r}) = \sum \rho_{G} \cdot \exp(i \cdot \underline{G} \cdot \underline{r})
$$

The vector *G* so far is just a mathematical construct defining the "inverse" space needed for the Fourier transform.

- However, since we can always substitute for any *r* a vector $r + T(T=$ translation vector of the lattice), or written out, $r + n_1 a_1 + n_2 a_2 + n_3 a_3$ with n_i = integers and a_i = base vectors of the lattice defining the crystal, the product $r \cdot G$ must not change its value if we substitute $r \cdot r \cdot n_1 a_1 + n_2 a_2 + n_3 a_3$.
- This requires that *G* **·** *T***=2π ·** *m* with *m***=** integer.
- This is essentially a definition of the vectors *G* that serve as the Fourier transforms of the vector *T*, i.e. the lattice in space. These *reciprocal lattice vectors*, as they are called, can be obtained from the base vectors defining the regular lattice in the following way:

If we write *G* in components we obtain

 $G = h \cdot g_1 + k \cdot g_2 + l \cdot g_3$

With *h, k, l*= integers.

The vectors *g***1**, *g***2**, and *g* **³** are then the *unit vectors* of the **reciprocal lattice**. *(yes – they are underlined, you just don't see it with some fonts!)*

If we now form the inner product of *G* **·** *T*, e.g., for simplicity, with *T***=** *n***1 ·** *a***1**, we obtain

$$
(h \cdot g_{1} + k \cdot g_{2} + l \cdot g_{3}) \cdot (n_{1} \cdot g_{1}) = 2\pi \cdot m
$$

For an arbitrary n_1 this only holds if

$$
g_1 \cdot \underline{a}_1 = 2\pi
$$

 $g_2 \cdot \underline{a}_1 = g_3 \cdot \underline{a}_1 = 0$

In general terms, we have

$$
g_i \cdot \underline{a}_j = 2 \pi \delta_{ij}
$$

With **δij** = **Kronecker** symbol, defined by: **δ ij=0** for **≠** and **δij =1** for **i=j**.

The above equation is satisfied with the following definitions for the unit vectors of the reciprocal lattice:

$$
g_1 = 2 \pi \frac{\underline{a_2} \times \underline{a_3}}{\underline{a_1} \cdot \underline{a_2} \cdot \underline{a_3}}
$$
\n
$$
g_2 = 2 \pi \frac{\underline{a_1} \cdot \underline{a_2} \cdot \underline{a_3}}{\underline{a_1} \cdot \underline{a_2} \cdot \underline{a_3}}
$$
\n
$$
g_3 = 2 \pi \frac{\underline{a_1} \times \underline{a_2}}{\underline{a_1} \cdot \underline{a_2} \cdot \underline{a_3}}
$$