## **Forward Currents from the Space Charge Region**

Again, we start from the equation for the net recombination  $U_{DL}$  via deep levels

$$
U_{\text{DL}} = \frac{v \cdot \sigma^{\text{e}} \cdot N_{\text{DL}} \cdot (n^{\text{e}} \cdot n^{\text{h}} - n_{\text{i}}^2)}{E_{\text{DL}} - E_{\text{MB}}}
$$
\n
$$
= \frac{E_{\text{DL}} - E_{\text{MB}}}{n^{\text{e}} + n^{\text{h}} + 2n_{\text{i}} \cdot \cosh \frac{E_{\text{DL}} - E_{\text{MB}}}{kT}}
$$
\n
$$
= \frac{E_{\text{DL}} - E_{\text{MB}}}{n^{\text{e}} + n^{\text{h}} + 2n_{\text{i}} \cdot \cosh \frac{E_{\text{DL}} - E_{\text{MB}}}{kT}}
$$

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with  $1/\tau = v \cdot \sigma^e \cdot N_{DL}$  as we know by now.

The carrier densities *n* **e** and *n* **<sup>h</sup>** [may be expressed via their Quasi-Fermi energies as](http://www.tf.uni-kiel.de/matwis/amat/semi_en/kap_2/backbone/r2_3_2.html#_2) *E* **Fe** and *E***<sup>F</sup> h** , respectively. For their product we get

$$
n^e \cdot n^h = n_1^2 \cdot \exp - \frac{E_F^e - E_F^h}{kT}
$$

For the forward direction we have  $E_F$ <sup>e</sup> –  $E_F$ <sup>h</sup> < 1 and thus

$$
n^{\rm e}, n^{\rm h} \gg n_{\rm i}
$$

This leaves us with

$$
U_{\text{DL}} = -\frac{1}{\tau} \frac{n^{\text{e}} \cdot n^{\text{h}}}{n^{\text{e}} + n^{\text{h}}}
$$

The *maximum* value for  $U_{DL}$  gives the upper limit for the net recombination rate and thus the *maximum current due to recombination in the SRC*, too. The maximum is defined by

$$
\frac{\partial \{(n^{\mathbf{e}} \cdot n^{\mathbf{h}}) / (n^{\mathbf{e}} + n^{\mathbf{h}})\}}{\partial n^{\mathbf{e}}} = \frac{\partial \{(n^{\mathbf{e}} \cdot n^{\mathbf{h}}) / (n^{\mathbf{e}} + n^{\mathbf{h}})\}}{\partial n^{\mathbf{h}}} = 0
$$

which gives us  $n^e = n^h$  for maximum current. With  $n^e \cdot n^h = n_i^2 \cdot \exp - [(E_F^e - E_F^h)]/k$   $T$  from above, we have

$$
n^{e} = n^{h} = n_{i} \cdot exp - \frac{E_{F}^{e} - E_{F}^{h}}{2kT}
$$

What we need now is an equation for the difference of the Quasi-Fermi energies. Lets look at the situation in a banddiagram



Whatever the exact positions of the Quasi-Fermi energies, their difference *E* **Fe –** *E***<sup>F</sup> h** is *about equal* to the difference in the bulk Fermi energy and thus

$$
E_F^e - E_F^h \approx e \cdot U
$$

(The "*about equal*" contains roughly the same approximation as the "[average barrier height"](http://www.tf.uni-kiel.de/matwis/amat/semi_en/kap_2/backbone/r2_2_4.html#_11) from the simple derivation!)

This gives us the final result

$$
U_{\text{DL}}\text{ (max)} \approx \frac{1}{2\tau} \cdot n_{\text{i}} \cdot \text{exp} - \frac{e \cdot U}{2kT}
$$

Again , this is the *net* recombination rate at any point in the space charge region. To obtain the current density, we have to multiply with the width *d* of the **SCR** (and the elementary charge) and obtain for the maximum current from the **SCR** in forward direction:

$$
j_{F}(SRC) = \frac{e \cdot n_{i} \cdot d_{SCR}}{2\tau} \cdot exp - \frac{e \cdot U}{2kT}
$$

Considering that we needed the whole formalism of Shockley-Read-Hall recombination theory, Quasi-Fermi energies, some junction theory, and lots of assumptions and approximations [to get the same result as before](http://www.tf.uni-kiel.de/matwis/amat/semi_en/kap_2/backbone/r2_2_4.html#_10), this does not appear to be a much better way of getting an idea about the influence of the **SCR** on the diode characteristic than the "quick and dirty" way.

But don't deceive yourself! The treatment given here is not only physically sound, but transparent at every step. If you want to do more precise calculations, you would know - at least in principle - what to do.